

## A Variational Iteration Method Involving Adomian Polynomials for a Strongly Nonlinear Boundary Value Problem

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**Abstract.** A variational iteration method involving Adomian polynomials to solve a strongly nonlinear boundary value problem is considered. After its convergence is established, the efficiency and accuracy of the proposed method are tested on problems with exponential nonlinearity.

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**Key words:** Boundary value problem, variational iteration method, Adomian polynomials, convergence.

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### 1. Introduction

The nonlinear boundary value problem

$$\begin{aligned}y'' + \frac{m}{x}y' + f(x, y) &= 0, \quad 0 \leq x \leq 1, \\ \alpha_1 y(0) + \beta_1 y'(0) &= \gamma_1, \quad m = 0 \quad \text{or} \quad y'(0) = 0, \quad m > 0, \\ \alpha_2 y(1) + \beta_2 y'(1) &= \gamma_2,\end{aligned}\tag{1.1}$$

where  $\alpha_i, \beta_i, \gamma_i, i = 1, 2$  are finite constants, arises in various applications, including thermal explosions [10], tumor growth models [7], electroosmotic flows [11, 12, 15], modelling of heat sources in human head [26], and oxygen diffusion [42]. The unique solvability of the problem (1.1) for  $m \geq 1$  and boundary conditions  $y'(0) = 0$  and  $y(1) = B$  was established by Chawla & Shivkumar [19], while the more general case of nonlinear boundary conditions was studied by Garner & Shivaji [27]. In order to find approximate solutions of the problem, various numerical methods have been used — e.g. the Adomian decomposition method [39], the Taylor series method [18], a variational iteration method [49, 58]. The first two of these methods experience convergence difficulties, while the third one is

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restricted to the solution of the problem (1.1) with functions  $f(x, y)$  belonging to a special class of non-linear polynomials. To avoid these difficulties, modifications of the last two methods have been suggested [13, 14, 16].

The variational iteration method proposed by He [30–32] and its modifications [38, 56] have a high convergence rate and small error, so they are widely regarded as a good tool for solving functional equations [1, 28, 29, 33, 34, 37, 41, 44, 45, 51, 53, 59] arising in nonlinear science and engineering problems [5, 6, 40, 50, 54, 55, 60, 61]. Based on the variational iteration algorithm I, the convergence of the method has been extensively studied in [29, 45, 51–53]. Recently, Chang [17] used the variational iteration algorithm II to prove the convergence of the method for two-point diffusion problems. For more details about the method and its applications we refer the reader to [36, 37] and references therein.

The concept of Adomian polynomials was introduced by Adomian [8] in 1976. Later, Adomian and Rach [9] presented a formal formula to generate the Adomian polynomials for all form of nonlinearity. Since then, various algorithms for calculating the Adomian polynomials have been proposed to improve computational efficiency [2, 22, 23, 47, 57]. Symbolic implementation of several recurrence algorithms by using *MATHEMATICA* or *MAPLE* was also developed — cf. [20, 21, 23, 46]. The convergence of the Adomian polynomial series has been also discussed in [2, 22, 48].

Here, we combine a variational iteration method with Adomian polynomials to obtain the approximate solutions of a strongly nonlinear boundary value problem. In contrast to the above mentioned methods, our approach does not require any additional tools. The key idea is that the nonlinear terms in the correction functionals are decomposed into a series of Adomian polynomials, as this simplifies the computations considerably. Sufficient conditions for the method convergence are established, and test examples involving exponential nonlinearity, demonstrate the efficiency of the algorithm. The errors do not depend on a specific location of a point  $x \in [0, 1]$ , and are valid for the whole domain considered. This approach can be applied to various problems involving other differential equations with a strong nonlinearity.

## 2. Convergence

According to He [35, 36] and Chang [17], the variational iteration algorithms I and II for the problem (1.1) have the form

$$\begin{aligned} y_{n+1}(x) &= y_n(x) + \int_0^x \lambda(s; x) \left[ y_n''(s) + \frac{m}{s} y_n'(s) + f(s, y_n(s)) \right] ds, \quad n \geq 0, \\ y_{n+1}(x) &= y_0(x) + \int_0^x \lambda(s; x) f(s, y_n(s)) ds, \quad n \geq 0, \end{aligned} \tag{2.1}$$

where  $y_0(x) := y(0) + y'(0)x$  and  $\lambda(s; x)$  is the Lagrange multiplier [49, 58] defined by

$$\lambda(s; x) = \begin{cases} s \ln\left(\frac{s}{x}\right), & m = 1, \\ \frac{s(s^{m-1} - x^{m-1})}{(m-1)x^{m-1}}, & 0 \leq m \neq 1. \end{cases} \quad (2.2)$$

Integrating the differential equation in (1.1) twice, one obtains

$$y'(x) = \begin{cases} y'(0) - \int_0^x f(s, y(s)) ds, & m = 0, \\ - \int_0^x \frac{s^m}{x^m} f(s, y(s)) ds, & m > 0, \end{cases} \quad (2.3)$$

$$y(x) = y_0(x) + \int_0^x \lambda(s; x) f(s, y(s)) ds. \quad (2.4)$$

Though the variational iteration method is widely used in nonlinear problems, the presence of complicated functions in the kernels of symbolic integrals can cause substantial difficulties in implementation. Therefore, we represent the nonlinear term  $f(x, y(x))$  as the series

$$f(x, y(x)) = \sum_{k=0}^{\infty} A_k(u_0(x), u_1(x), \dots, u_k(x)), \quad (2.5)$$

where  $u_0(x) = y_0(x)$ ,  $u_k(x) = y_k(x) - y_{k-1}(x)$ ,  $A_0(u_0(x)) = f(x, u_0(x))$ , and for  $k \geq 1$  the Adomian polynomials  $A_k(x)$  are recurrently defined in [23] by

$$A_k(x) = A_k(u_0(x), u_1(x), \dots, u_k(x)) = \sum_{i=1}^k f^{(i)}(x, u_0(x)) C_k^i(x), \quad k \geq 1 \quad (2.6)$$

with

$$C_k^1(x) := u_k(x), \quad k \geq 1, \quad C_k^i(x) := \frac{1}{k} \sum_{j=0}^{k-i} (j+1) u_{j+1}(x) C_{k-1-j}^{i-1}(x), \quad 2 \leq i \leq k.$$

Such a modification is called the variational iteration-Adomian method [3, 4, 24, 25]. Analogously, taking into account the Eq. (2.4), one obtains

$$y(x) = y_0(x) + \int_0^x \lambda(s; x) \sum_{k=0}^{\infty} A_k(s) ds \quad (2.7)$$

and for  $n \geq 0$  the variational iteration algorithm I and algorithm II take the form

$$y_{n+1}(x) = y_n(x) + \int_0^x \lambda(s; x) \left[ y_n''(s) + \frac{m}{s} y_n'(s) + \sum_{k=0}^n A_k(s) \right] ds, \quad (2.8)$$

$$y_{n+1}(x) = y_0(x) + \int_0^x \lambda(s; x) \sum_{k=0}^n A_k(s) ds. \quad (2.9)$$

The Adomian polynomials  $A_k(x)$  above are also called the classical Adomian polynomials. According to [48], they satisfy the relations

$$\begin{aligned} A_0(u_0(x)) &= f(x, u_0(x)) = f(x, y_0(x)), \\ A_k(x) &< f(x, y_k(x)) - f(x, y_{k-1}(x)), \quad k \geq 1. \end{aligned} \quad (2.10)$$

Following Ref. [17], we present the truncation error for the method under consideration and establish sufficient conditions for its convergence.

**Theorem 2.1.** *Assume that the function  $f(x, y)$  is analytic in the rectangle  $R = \{(x, y) : 0 \leq x \leq 1, |y - y_0| \leq b\}$  and bounded from above by an  $M$  such that  $M < 2(m+1)b$ . Then every sequence (2.8) or (2.9) with  $y_0(x) = y(0) + y'(0)x$  converges to the exact solution  $y(x)$  of the problem (1.1) in the norm of  $C_{[0,1]}$ . Moreover,*

$$E_n := \|y(x) - y_n(x)\|_{C_{[0,1]}} \leq \frac{MK^n}{(n+1)!2^{n+1}(m+1)(m+3)\cdots(m+2n+1)}, \quad (2.11)$$

where  $K = \|\partial f(x, y)/\partial y\|_{C_{[0,1]}}$ .

*Proof.* Since  $f(x, y)$  is analytic in  $R$ , by the Cauchy-Kovalevskaya theorem the problem (1.1) has unique solution  $y(x)$  analytic in  $x \in [0, 1]$ . Therefore, the Murray-Miller theorem for existence and uniqueness [43] guarantees that  $f(x, y) = \sum_{k=0}^{\infty} A_k(u_0, u_1, \dots, u_k)$ , where  $A_k(x)$  is the classical Adomian polynomials, and the above series is uniformly convergent because it is essentially a rearrangement of the parameterised Taylor expansion series of the analytic function  $f(x, y)$  about the function  $u_0(x)$  [48]. Obviously,  $(x, y_0) \in R$  for all  $x \in [0, 1]$  such that  $|f(x, y_0)| = |A_0(x)| \leq M$ . It follows from (2.9) and (A.4) that for  $n = 0$  one has

$$|y_1(x) - y_0(x)| \leq \int_0^x |\lambda(s; x)| |A_0(s)| ds \leq \frac{M}{2(m+1)} x^2 \leq \frac{M}{2(m+1)} < b,$$

so that  $(x, y_1) \in R$  and  $|f(x, y_1)| \leq M$ . Further, the usual induction arguments show that  $(x, y_n) \in R$  and  $|f(x, y_n)| \leq M$  for any positive integer  $n$ . Thus the expressions containing the points  $y_n(x)$ ,  $n \in \mathbb{N}$  can be estimated — e.g. the Eqs. (2.7) and (A.4) yield

$$\begin{aligned} |y(x) - y_0(x)| &\leq \int_0^x |\lambda(s; x)| \left| \sum_{k=0}^{\infty} A_k(u_0, u_1, \dots, u_k) \right| ds = \int_0^x |\lambda(s; x)| |f(s, y(s))| ds \\ &\leq M \int_0^x |\lambda(s; x)| ds \leq \frac{M}{2(m+1)} x^2 \leq \frac{M}{2(m+1)}. \end{aligned}$$

Since  $f(x, y)$  is an analytic function in  $R$  and  $(x, y_n) \in R$ , there is a constant  $K = \|\partial f(x, y)/\partial y\|_{C_{[0,1]}}$  such that

$$|f(x, y) - f(x, y_n)| \leq K |y(x) - y_n(x)|, \quad x \in [0, 1].$$

The Eq. (2.10) implies that

$$\left| \sum_{k=n+1}^{\infty} A_k(u_0, u_1, \dots, u_k) \right| \leq K |y(x) - y_n(x)|, \quad x \in [0, 1].$$

Analogously, if  $x \in [0, 1]$ , then

$$\begin{aligned} |y(x) - y_1(x)| &\leq \int_0^x |\lambda(s; x)| \left| \sum_{k=1}^{\infty} A_k(u_0, u_1, \dots, u_k) \right| ds \\ &\leq K \int_0^x |\lambda(s; x)| |y(s) - y_0(s)| ds \\ &\leq \frac{MK}{2(m+1)} \int_0^x |\lambda(s, x) s^2| ds \leq \frac{MK}{2!2^2(m+1)(m+3)} x^4 \\ &\leq \frac{MK}{2!2^2(m+1)(m+3)}. \end{aligned}$$

Continuing in this way, we eventually arrive at the estimate

$$|y(x) - y_n(x)| \leq \frac{MK^n}{(n+1)!2^{n+1}(m+1)(m+3) \cdots (m+2n+1)}$$

valid for all  $x \in [0, 1]$ . Therefore, the inequality (2.11) holds and  $E_n \rightarrow 0$  as  $n \rightarrow \infty$ .  $\square$

### 3. Applications

In this section, we consider approximate solution of several boundary value problems (1.1) found with the assistance of *MATHEMATICA 5.2* on a laptop with a Pentium M 1.4GHz and 256 MB of RAM.

**Example 3.1.** The nonlinear singular boundary value problem

$$\begin{aligned} y'' + \frac{1}{x}y' + e^y &= 0, \\ y'(0) = 0, \quad y(1) &= 0 \end{aligned} \tag{3.1}$$

arises in the theory of thermal explosions [10].

For  $m = 1$ , the correction functional has the form

$$y_{n+1}(x) = y_n(x) + \int_{x_0}^x s \ln\left(\frac{s}{x}\right) \left[ y_n''(s) + \frac{1}{s}y_n'(s) + e^{y(s)} \right] ds.$$

Choosing a constant  $a$  as the initial approximation  $y_0$ , we obtain

$$y_1(x) = a - \frac{e^a}{4}x^2.$$

However, the second and further consecutive iterations cannot be computed in a closed form due to complicated integrands. This is a common complication in the implementation of the variational iteration method for strongly nonlinear problems. In contrast, in the variational iteration-Adomian method, the term  $y_{n+1}(x)$  is calculated by the formula

$$y_{n+1}(x) = y_n(x) + \int_0^x s \ln\left(\frac{s}{x}\right) \left[ y_n''(s) + \frac{1}{s} y_n'(s) + \sum_{k=0}^n A_k(s) \right] ds.$$

According to the algorithm (2.6), the first four Adomian polynomials are

$$\begin{aligned} A_0 &= e^{u_0}, & A_1 &= e^{u_0} u_1, \\ A_2 &= \frac{1}{2} e^{u_0} u_1^2 + e^{u_0} u_2, \\ A_3 &= \frac{1}{6} e^{u_0} u_1^3 + e^{u_0} u_1 u_2 + e^{u_0} u_3 \end{aligned}$$

produce the following approximations of the solution  $y$ :

$$\begin{aligned} y_1(x) &= a - \frac{e^a}{4} x^2, \\ y_2(x) &= a - \frac{e^a}{4} x^2 + \frac{e^{2a}}{64} x^4, \\ y_3(x) &= a - \frac{e^a}{4} x^2 + \frac{e^{2a}}{64} x^4 - \frac{e^{3a}}{768} x^6, \\ y_4(x) &= a - \frac{e^a}{4} x^2 + \frac{e^{2a}}{64} x^4 - \frac{e^{3a}}{768} x^6 - \frac{e^{4a}}{8192} x^8. \end{aligned}$$

It follows from the Eq. (2.3) that  $y'(x) \leq 0$  for all  $x \in [0, 1]$ . Hence,  $y(x)$  is a non-increasing function ( $0 \leq y(x) \leq a$  for  $0 \leq x \leq 1$ ), which implies  $M = K = e^a$ . Thus the sequence  $y_n(x)$  converges to the exact solution of the problem (3.1) in the space  $C_{[0,1]}$  and

$$E_n \leq \frac{e^{a(n+1)}}{2^{2(n+1)}(n+1)!(n+1)!}. \tag{3.2}$$

The unknown constant  $a$  is to be determined by imposing a boundary condition on  $y_n$  at  $x = 1$  to obtain a transcendental equation,  $y_n(1) = 0$ . *MATHEMATICA* has a built-in command to solve this transcendental equation. Once the unknown constant  $a$  has been determined, the corresponding maximum absolute error is obtained from the Eq. (3.2). A sequence of approximations for  $a$  and the corresponding maximum absolute errors for  $n = 5, 10$  and  $12$  are shown in Table 1. As expected, this sequence converges to  $a$  with the high convergence rate.

Table 1: Example 3.1. Constant  $a$  and maximum absolute errors (MAEs).

$n$	5	10	12	Exact [39]
$a$	0.3167048552978	0.316694366798	0.3166943676198	0.3166943676407
MAE	$3.15 \times 10^{-9}$	$4.88 \times 10^{-21}$	$2.36 \times 10^{-26}$	–

**Example 3.2.** The nonlinear singular boundary value problem

$$\begin{aligned} y'' + \frac{2}{x}y' + e^{-y} &= 0, \\ y'(0) = 0, \quad 0.1y(1) + y'(1) &= 0 \end{aligned} \tag{3.3}$$

is used in a heat conduction model of the human head [26].

In this case, the variational iteration algorithm II is

$$y_{n+1}(x) = y_0(x) + \int_0^x \frac{s(s-x)}{x} \sum_{k=0}^n A_k(s) ds,$$

and the corresponding first four Adomian polynomials are

$$\begin{aligned} A_0 &= e^{-u_0}, \quad A_1 = -e^{-u_0}u_1, \\ A_2 &= \frac{1}{2}e^{-u_0}u_1^2 - e^{-u_0}u_2, \\ A_3 &= -\frac{1}{6}e^{-u_0}u_1^3 + e^{-u_0}u_1u_2 - e^{-u_0}u_3. \end{aligned}$$

Taking  $y_0 = a$  as the initial approximation, produces

$$\begin{aligned} y_1(x) &= a - \frac{e^{-a}}{6}x^2, \\ y_2(x) &= a - \frac{e^{-a}}{6}x^2 - \frac{e^{-2a}}{120}x^4, \\ y_3(x) &= a - \frac{e^{-a}}{6}x^2 - \frac{e^{-2a}}{120}x^4 - \frac{e^{-3a}}{1890}x^6, \\ y_4(x) &= a - \frac{e^{-a}}{6}x^2 - \frac{e^{-2a}}{120}x^4 - \frac{e^{-3a}}{1890}x^6 - \frac{61e^{-4a}}{1632960}x^8. \end{aligned}$$

In this case,  $M = K = 1$  — cf. Ref. [17]. Consequently, the sequence  $y_n(x)$  converges to the exact solution in the space  $C_{[0,1]}$  and

$$E_n \leq \frac{1}{(2n+3)!}.$$

The approximations of  $a$  and the corresponding maximum absolute errors are presented in Table 2. It shows the rapid convergence of the approximation sequence.

Table 2: Example 3.2. Constant  $a$  and maximum absolute errors (MAEs).

$n$	1	5	9	10
$a$	1.130289326974	1.147039006534	1.14703901933	1.14703901933
MAE	$8.34 \times 10^{-3}$	$1.61 \times 10^{-10}$	$1.96 \times 10^{-20}$	$3.87 \times 10^{-23}$

**Example 3.3.** Consider the following nonlinear boundary value problem:

$$\begin{aligned} y'' + e^{x+y} &= 1, \\ y'(0) &= -1, \quad y(1) = -1. \end{aligned} \quad (3.4)$$

The corresponding variational iteration algorithm II has the form

$$y_{n+1}(x) = y_0(x) + \int_0^x (s-x) \sum_{k=0}^n A_k(s) ds,$$

and the first four Adomian polynomials are

$$\begin{aligned} A_0 &= e^{x+u_0} - 1, \quad A_1 = e^{x+u_0} u_1, \\ A_2 &= \frac{1}{2} e^{x+u_0} u_1^2 + e^{x+u_0} u_2, \\ A_3 &= \frac{1}{6} e^{x+u_0} u_1^3 + e^{x+u_0} u_1 u_2 + e^{x+u_0} u_3. \end{aligned}$$

Taking  $y_0 = a - x$  as the initial approximation, we obtain

$$\begin{aligned} y_1(x) &= a - x - \frac{e^a - 1}{2} x^2, \\ y_2(x) &= a - x - \frac{e^a - 1}{2} x^2 + \frac{e^a(e^a - 1)}{24} x^4, \\ y_3(x) &= a - x - \frac{e^a - 1}{2} x^2 + \frac{e^a(e^a - 1)}{24} x^4 - \frac{3e^a - 7e^{2a} + 4e^{3a}}{720} x^6, \\ y_4(x) &= a - x - \frac{e^a - 1}{2} x^2 + \frac{e^a(e^a - 1)}{24} x^4 - \frac{3e^a - 7e^{2a} + 4e^{3a}}{720} x^6 \\ &\quad - \frac{15e^a - 63e^{2a} + 82e^{3a} - 34e^{4a}}{40320}. \end{aligned}$$

Arguments analogous to those in Example 3.1 show that  $M = e^{a+1} - 1$ ,  $K = e^{a+1}$  and

$$E_n \leq \frac{e^{n(a+1)}(e^{a+1} - 1)}{(2n+2)!}.$$

Proceeding as before, the values of  $a$  for  $n = 1, 10$  and  $20$  are found to be the same as  $0$ . This leads to function  $y(x) = -x$ , which is the exact solution of the problem.

#### 4. Conclusion

We apply a variational iteration method using Adomian polynomials to strongly nonlinear boundary value problems. This approach involves the calculation of symbolic integrals connected with variational iteration methods. Error estimates have been derived and the convergence of the method established. Illustrative examples demonstrate the efficiency of the algorithm.



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### A. Appendix

Define a function  $g(x)$  by

$$g(x) = \int_0^x \lambda(s; x) s^k ds, \quad k \geq 0, \quad (\text{A.1})$$

where  $\lambda(s; x)$  is the Lagrange multiplier (2.2) such that

$$\begin{aligned} \lambda(s = x; x) &= 0, \\ \left. \frac{\partial \lambda(s; x)}{\partial x} \right|_{s=x} &= -1, \\ \frac{\partial^2 \lambda(s; x)}{\partial x^2} + \frac{m}{x} \frac{\partial \lambda(s; x)}{\partial x} &= 0. \end{aligned} \quad (\text{A.2})$$

Calculating the derivative of  $g$  in  $x$  and using Leibnitz formula yields

$$\begin{aligned} g'(x) &= \int_0^x \frac{\partial \lambda(s; x)}{\partial x} s^k ds, \\ g''(x) &= \int_0^x \frac{\partial^2 \lambda(s; x)}{\partial x^2} s^k ds - x^k. \end{aligned} \quad (\text{A.3})$$

It follows that

$$g''(x) + \frac{m}{x} g'(x) + x^k = 0,$$

or

$$\frac{d}{dx} [x^m g'(x)] = -x^{k+m}.$$

Integrating the above expression twice and using the conditions  $g(0) = g'(0) = 0$ , one obtains

$$g(x) = \int_0^x \lambda(s; x) s^k ds = \frac{-1}{(k+m+1)(k+2)} x^{k+2}. \quad (\text{A.4})$$

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