

## A HYPERBOLIC-ELLIPTIC MODEL OF TWO-PHASE FLOW IN POROUS MEDIA – EXISTENCE OF ENTROPY SOLUTIONS

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*Dedicated to the memory of Magne S. Espedal*

**Abstract.** We consider the flow of two-phases in a porous medium and propose a modified version of the fractional flow model for incompressible, two-phase flow based on a Helmholtz regularization of the Darcy phase velocities. We show the existence of global-in-time entropy solutions for this model with suitable assumptions on the boundary conditions. Numerical experiments demonstrating the approximation of the classical two-phase flow equations with the new model are presented.

**Key Words.** Porous media flow, conservation law, elliptic equation, weak solution, existence

### 1. The two Phase Flow Problem

Many geophysical and industrial processes like enhanced oil recovery and carbon dioxide sequestration involve the flow of two-phases, say oil and water, in a porous medium.

The variables of interest are the phase saturations  $s_w$  and  $s_o$  representing the saturation (volume fraction) of the water and oil phase respectively. We have the identity:

$$(1.1) \quad s_w + s_o \equiv 1.$$

Hence, we can describe the dynamics in terms of the saturation of either of the two-phases. We denote the water saturation as  $s_w = s$  in the discussion below. Assuming a constant porosity ( $\phi \equiv 1$ ), the two-phases are transported by [4]

$$(1.2) \quad (s_r)_t + \operatorname{div}_x(\mathbf{v}_r) = 0, \quad r \in \{w, o\}.$$

Here, the phase velocities are denoted by  $\mathbf{v}_w$  and  $\mathbf{v}_o$  respectively. In view of the identity (1.1), the two-phase velocities can be summed up to yield the *incompressibility* condition,

$$(1.3) \quad \operatorname{div}_x(\mathbf{v}) = 0, \quad \mathbf{v} = \mathbf{v}_w + \mathbf{v}_o.$$

The total velocity is denoted by  $\mathbf{v}$ .

The phase velocities in a homogeneous isotropic medium are described by the Darcy's law [4]:

$$(1.4) \quad \mathbf{v}_r = -\lambda_r \nabla_x p_r + \lambda_r \rho_r g \mathbf{k}, \quad r \in \{w, o\}.$$

Here,  $g$  is the constant acceleration due to gravity,  $\mathbf{k}$  is the direction in which gravity acts and  $\rho_r$  is the (constant) density of the phase  $r$ . The quantity  $\lambda_r = \lambda_r(s_r)$  is

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the phase mobility and  $p_r$  is the phase pressure. Assume that the *capillary pressure* i.e,  $p_c = p_w - p_o$  is zero, we can sum (1.4) for both phases and obtain

$$(1.5) \quad \mathbf{v} = -\lambda_T(s)\nabla_x p + (\lambda_w\rho_w + \lambda_o\rho_o)g\mathbf{k},$$

with  $p = p_w = p_o$  being the pressure and  $\lambda_T = \lambda_w + \lambda_o$  being the total mobility. Using (1.5), the gradient of pressure in (1.4) can be eliminated leading to

$$\mathbf{v}_w = \frac{\lambda_w(s)}{\lambda_T(s)}\mathbf{v} + \frac{\lambda_w(s)\lambda_o(s)}{\lambda_T(s)}(\rho_w - \rho_o)g\mathbf{k}.$$

Denoting the fractional flow function  $f$  as

$$f(s) = \frac{\lambda_w(s)}{\lambda_T(s)} = \frac{\lambda_w(s)}{\lambda_w(s) + \lambda_o(s)},$$

and the gravity function  $g$  as

$$g(s) = \frac{\lambda_w(s)\lambda_o(s)}{\lambda_T(s)}(\rho_w - \rho_o)g,$$

the saturation equation (1.2) for water can be written down as

$$(1.6) \quad s_t + \operatorname{div}_x(f(s)\mathbf{v} + g(s)\mathbf{k}) = 0.$$

Combining the saturation equation with the incompressibility condition (1.3) and the pressure equation, we obtain the evolution equations for two-phase flow in a porous medium:

$$(1.7) \quad \begin{aligned} s_t + \operatorname{div}_x(f(s)\mathbf{v} + g(s)\mathbf{k}) &= 0, \\ \operatorname{div}_x(\mathbf{v}) &= 0, \\ \mathbf{v} &= -\lambda_T(s)\nabla_x p + (\rho_w\lambda_w(s) + \rho_o\lambda_o(s))g\mathbf{k}. \end{aligned}$$

The above equations have to be augmented by suitable initial and boundary conditions.

The phase mobility  $\lambda_w : [0, 1] \mapsto \mathbb{R}$  is a monotone increasing function with  $\lambda_w(0) = 0$  and the phase mobility  $\lambda_o : [0, 1] \mapsto \mathbb{R}$  is a monotone decreasing function with  $\lambda_o(1) = 0$ . Furthermore, the total mobility is strictly positive i.e,  $\lambda_T \geq \lambda_* > 0$  for some  $\lambda_*$ .

The above equations are a hyperbolic-elliptic system as the saturation equation in (1.7) is a scalar hyperbolic conservation law in several space dimensions with a coefficient given by the velocity  $\mathbf{v}$ . The velocity can be obtained by solving an elliptic equation for the pressure  $p$ .

It is well known that solutions of hyperbolic conservation laws can develop discontinuities, even for smooth initial data, [8]. The presence of these discontinuities or shock waves implies that solutions of conservation laws are sought in a weak sense and are augmented with additional admissibility criteria or *entropy conditions* in order to ensure uniqueness.

As the two-phase flow equations involve a conservation law, we need to define a suitable concept of entropy solutions for these equations and show that these solutions are well-posed. The problem of proving well-posedness of global weak solutions of the two-phase flow equations (1.7) has remained open for many decades. The main challenge in showing existence is the fact that the velocity field  $\mathbf{v}$  acts as a coefficient in the saturation equations. Although conservation laws with coefficients have been studied extensively in recent years, see [1, 11, 7, 2] and references therein, the state of the art results require that the coefficient is a function of bounded variation. Many attempts at showing that the velocity field  $\mathbf{v}$  in (1.6) is sufficiently regular, for example is a *BV* function or has enough Sobolev regularity, have failed.