

An algebraic expression of the three-dimensional Franck-Condon factors and its application

Hong-Mo Huang^{a,*} and Jun Liang^{b,*}

^a Department of Physics, Anhui Tongcheng Normal College, Tongcheng 231400, China

^b College of Physics and Electronic Information, Anhui Normal University, Wuhu 241000, China

Received 1 December 2010; Accepted (in revised version) 14 January 2011

Published Online 28 February 2011

Abstract. A more general algebraic expression for the calculation of the four-mode Franck-Condon factors was derived straightforwardly on the base of the closed form expression of the Franck-Condon integrals between arbitrary multidimensional harmonic oscillators under the Duschinsky mixing effects. This new algebraic expression was applied to study the photoelectron spectra of $D_2CO^+(\tilde{A}^2B_1)$. Franck-Condon analyses and spectral simulations were carried out on the $D_2CO^+(\tilde{A}^2B_1) - D_2CO(\tilde{X}^1A_1)$ photoionization processes. The spectral simulations of vibrational structures based on the computed Franck-Condon factors are in excellent agreement with the observed spectra.

PACS: 31.15.-p, 31.15.xr, 33.15.-e

Key words: overlap integral, Franck-Condon factor, duschinsky effect, spectral simulation

1 Introduction

The square of the vibrational overlap integral between two electronic states is called the Franck-Condon factor (FCF). Calculations of FCFs are crucial for interpreting vibronic spectra of molecules as well as studying nonradiative processes. Recently, we have developed a new method for calculating Franck-Condon factors of multidimensional harmonic oscillators including the Duschinsky effect [1, 2]. Some explicit algebraic formulas of two-dimensional (two-, three-, and four-mode) Franck-Condon factors were derived straightforwardly by the properties of Hermite polynomials and Gaussian integrals. This new method was applied to study the photoelectron spectra of ClO_2^- , SO_2 , CH_3OO^- and so on [3–7]. Our approach is alternative to other existing ones [8–19] and has the advantages of being efficient and having

*Corresponding author. Email addresses: hhmd@163.com (H. M. Huang); jliang@mail.ahnu.edu.cn (J. Liang)

no singular points. Accordingly, our method can be applied to any displaced-distorted-rotated harmonic oscillators and should be valuable in the studies of vibronic spectroscopy and non-radiative processes of molecules. However, up to date, an explicit algebraic form expression to calculate the three-dimensional four-mode Franck-Condon factors under the Duschinsky mixing effects has not been reported according to our knowledge.

In this work, we extended our approach to calculate three-dimensional Franck-Condon factors. An analytical expression for the calculation of the three-dimensional four-mode Franck-Condon integrals has been exactly derived. In addition, a general explicit formula of the three-dimensional Franck-Condon factors was given. As an example we present a calculation of the intensity distribution in the photoelectron spectrum of the $D_2CO^+(\tilde{A}^2B_1) - D_2CO(\tilde{X}^1A_1)$ transition of Formaldehyde.

2 Theoretical method

In Refs. [1, 2], a closed form expression for multidimensional Franck-Condon integrals between displaced distorted-rotated harmonic potential surfaces has been derived

$$\begin{aligned} \langle v''_1 \cdots v''_n | v'_1 \cdots v'_N \rangle = & \langle 0''_1 \cdots 0''_N | 0'_1 \cdots 0'_N \rangle \left(\prod_{j=1}^N (-1)^{v''_j + v'_j} (v''_j! v'_j!)^{-1/2} \right) \\ & \times \exp \left(\frac{1}{2} \boldsymbol{\sigma}''^T (\mathbf{I} - 2\mathbf{Q}) \boldsymbol{\sigma}'' + \frac{1}{2} \boldsymbol{\sigma}'^T (\mathbf{I} - 2\mathbf{P}) \boldsymbol{\sigma}' - 2\boldsymbol{\sigma}''^T \mathbf{R} \boldsymbol{\sigma}' \right) \\ & \times \frac{\partial^{v''_1 + \cdots + v''_N + v'_1 + \cdots + v'_N}}{\partial \sigma_1''^{v''_1} \cdots \partial \sigma_N''^{v''_N} \partial \sigma_1'^{v'_1} \cdots \partial \sigma_N'^{v'_N}} \\ & \times \exp \left(-\frac{1}{2} \boldsymbol{\sigma}''^T (\mathbf{I} - 2\mathbf{Q}) \boldsymbol{\sigma}'' - \frac{1}{2} \boldsymbol{\sigma}'^T (\mathbf{I} - 2\mathbf{P}) \boldsymbol{\sigma}' + 2\boldsymbol{\sigma}''^T \mathbf{R} \boldsymbol{\sigma}' \right), \quad (1) \end{aligned}$$

where

$$\langle 0''_1 \cdots 0''_N | 0'_1 \cdots 0'_N \rangle = 2^{N/2} (\det \Gamma' \Gamma'')^{1/4} (\det \mathbf{JQ})^{1/2} \exp \left(-\frac{1}{2} \boldsymbol{\delta}^T (\mathbf{I} - \mathbf{P}) \boldsymbol{\delta} \right), \quad (2)$$

and

$$\begin{pmatrix} \boldsymbol{\sigma}'' \\ \boldsymbol{\sigma}' \end{pmatrix} = \sqrt{2} \begin{pmatrix} \mathbf{I} - 2\mathbf{Q} & -2\mathbf{R} \\ -2\mathbf{R}^T & \mathbf{I} - 2\mathbf{P} \end{pmatrix}^{-1} \begin{pmatrix} -\mathbf{R}\boldsymbol{\delta} \\ (\mathbf{I} - \mathbf{P})\boldsymbol{\delta} \end{pmatrix}. \quad (3)$$

Here \mathbf{I} is an $N \times N$ unit matrix, and symmetric matrices \mathbf{P} and \mathbf{Q} and the $N \times N$ matrix \mathbf{R} are defined by

$$\mathbf{P} = \mathbf{S}\mathbf{Q}\mathbf{S}^T, \quad \mathbf{Q} = (\mathbf{1} + \mathbf{S}^T \mathbf{S})^{-1}, \quad \mathbf{R} = \mathbf{Q}\mathbf{S}^T, \quad (4)$$

with

$$\mathbf{S} = \boldsymbol{\lambda}_\omega \mathbf{J} \boldsymbol{\lambda}_{\omega'}^{-1}, \quad \boldsymbol{\lambda}_\omega = \text{diag}(\sqrt{\omega_1}, \dots, \sqrt{\omega_N}), \quad \boldsymbol{\delta} = \hbar^{-1/2} \boldsymbol{\lambda}_{\omega'} \mathbf{K}. \quad (5)$$