

Influences of the low collision energies and isotopic variants on the stereodynamics for reactions $\text{Li}+\text{DF}/\text{TF}\rightarrow\text{LiF}+\text{D}/\text{T}$

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Abstract. Stereodynamics for $\text{Li}+\text{DF}/\text{TF}\rightarrow\text{LiF}+\text{D}/\text{T}$ reactions are studied using the quasi-classical trajectory method based on a new potential energy surface constructed by Aguado and Paniaga [J. Chem. Phys. 119 (2003) 10088]. The product angular distributions of $P(\theta_r)$ and $P(\phi_r)$, which reflect the vector correlation, are calculated and discussed. The average rotational alignment factor $\langle P_2(\cos\theta_r) \rangle$ as function of the collision energy is also presented. Furthermore, four polarization-dependent differential cross sections (PDDCSs), namely, PDDCS₀₀, PDDCS₂₀, PDDCS₂₁, and PDDCS₂₂₊ are calculated as well. By comparing the stereodynamics results of the title reactions, we find that the isotopic effects are relatively obvious.

Keywords: Quasi-classical trajectory; Stereodynamics; Product angular distribution; Polarization-dependent generalized differential cross sections; Mass factor

1. Introduction

The $\text{Li}+\text{HF}\rightarrow\text{LiF}+\text{H}$ reaction, which serves as a prototype for textbook reactions between alkali metals and hydrogen halides, has attracted a lot of attentions from both experimental[1-5] and theoretical[6-16] aspects. The first crossed molecular beam experiment on the LiFH system was carried out by Becker *et al.* in 1980 to detect the product angular distribution.[1] Subsequently, the $\text{Li}+\text{HF}(v=1, j=1, m=0)\rightarrow\text{LiF}+\text{H}$ reaction was studied in detail by Loesch *et al.* at the collision energy of 420 meV.[2-4] Using crossed molecular beam apparatus, Höbel *et al.* measured the double differential cross sections for the $\text{Li}+\text{HF}\rightarrow\text{LiF}+\text{H}$ reaction at collision energies ranging from 88 to 378 meV.[5]

In terms of theory, the $\text{Li}+\text{HF}\rightarrow\text{LiF}+\text{H}$ reaction, with few electrons, is relative simple, which provides likelihood to perform high-quality *ab initio* calculations for constructing an ideal potential energy surface (PES). There is a relatively deep van der Waals well in the reagent valley and a strongly bent transition state with a considerable barrier in the exit channel,[17] which makes the title reaction system possessing abundant dynamic information. In 1995, Aguado *et al.* constructed the PES of the reaction using the multiple reference single and double excitations configuration-interaction (MRDCI) method.[18] Soon afterwards, the same group presented a new fit to the more accurate MRDCI *ab initio* PES.[14] Moreover, they carried out the quantum dynamics study based on local coordinates using a three-dimensional time-dependent method, which makes it possible to describe reactants and products at the same time. In 2003, Aguado *et al.* performed high level *ab initio* calculations for a large number of nuclear configurations (about 6000) based on a new atomic basis sets and fitted four analytic global PESs for three $^2\text{A}'$ and one $^2\text{A}''$ electronic states. [17] Furthermore, they simulated spectrum of the LiHF system with the new PES, which is in very good agreement with the experimental

data of Hudson *et al.*[19]

Lagana *et al.* presented zero total angular momentum exact quantum probabilities for $\text{Li}+\text{HF}$ and its isotopic variants reactions in 2000.[20] Wecka *et al.* carried out a quantum reactive scattering calculations for zero total angular momentum at low and ultralow temperatures in 2005.[8] Based on the AP2-PES(1997), Yuan and Zhao explored the scalar and vector properties using the quasi-classical trajectory (QCT) method, aiming at studying the stereodynamics features of the reaction system in 2010.[21] Recently, based on the ground state PES, Liu *et al.* studied the vector properties of the $\text{Li}+\text{HF}\rightarrow\text{LiF}+\text{H}$ as well as its isotopic variant reactions at translational collision energies ranging from 30 kcal/mol to 60 kcal/mol using the QCT method.[22] Soon afterwards, the same group investigated the stereodynamics properties for the same reaction at collision energies ranging from 1.15 kcal/mol to 5.0 kcal/mol.[23] However, except the research carried out by Liu *et al.*, [22, 23] it is rare to study the LiHF system based on the new PES. Nevertheless, the Liu's group didn't investigate the isotopic substitution of the system when the collision energy is not more than 5.0 kcal/mol. In order to fully characterize stereodynamics of the isotopic substitution of $\text{Li}+\text{HF}$ reaction at lower collision energy, we explore the vector properties for $\text{Li}+\text{DF}/\text{TF}\rightarrow\text{LiF}+\text{D}/\text{T}$ reactions using the QCT method based on the new PES constructed by Aguado *et al.* over a collision energy range of 2.5-5.0 kcal/mol.

2. Theory

In this work, the adopted accurate full dimensional PES constructed by Aguado *et al.*[17] which takes the following form written as,

$$V_{ABC} = \sum_A V_A^{(1)} + \sum_{AB} V_{AB}^{(2)}(R_{AB}) + V_{ABC}^{(3)}(R_{AB}, R_{AC}, R_{BC})$$

Where $V_A^{(1)}$ is the energy of atom A in its appropriate electronic state, and the value of $\sum_A V_A^{(1)}$ is usually taken as zero because all the atoms are in their ground state. Moreover, $V_{AB}^{(2)}$ and $V_{ABC}^{(3)}$ are the two-body and three-body energy term respectively. In addition, R_{AB} , R_{AC} , and R_{BC} represent the distances of AB, AC, BC respectively.

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Based on the new PES, we carried out the dynamical calculations by employing the stereo-QCT procedure developed by Han et al.[24-31] The classical Hamilton equations are numerically integrated in three dimensions. In order to ensure the conservation of total energy and angular momentum, the integration step size is chosen to be 0.1 fs. The impact parameter b is optimized through the repeated computations with 100000 trajectories. The maximum impact parameter b_{max} is obtained where there is no reactive trajectory anymore when the value of b is slightly increased. The present QCT calculations are performed in the center-of-mass (CM) frame, depicted in Figure 2.1.

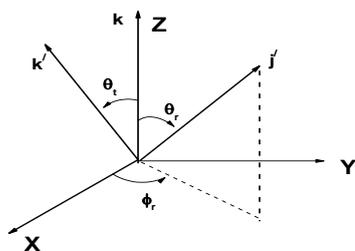


Figure 2.1: Centre-of-mass coordinate system used to describe the k , k' and j' correlations.

3. Results and discussion

Figure 3.1 displays the product angular distributions of $P(\theta_r)$ representing $k - j'$ correlation at different collision energies. As can be seen in Figure 3.1 (a), all $P(\theta_r)$ distributions of $\text{Li+DF} \rightarrow \text{LiF+D}$ reaction at different collision energies from 2.5 kcal/mol to 5.0 kcal/mol show a prominent peak at $\theta_r = 90^\circ$ and are symmetric with respect to $\theta_r = 90^\circ$, which means that the product rotational angular momentum vector j' is perpendicular to the direction of the reagent relative velocity k and indicates that the $P(\theta_r)$ distribution has a strong product rotational alignment. Furthermore, there is an obvious trend that the peaks of the product $P(\theta_r)$ distribution become higher with the increase of the collision energy from 2.5 kcal/mol to 5.0 kcal/mol, which shows that the product alignment becomes stronger as the collision energy increases. In addition, in Figure 3.1 (b), there exist almost the same phenomena except a little bit discrepancy for the isotopic reaction $\text{Li+TF} \rightarrow \text{LiF+T}$ at different collision energies from 2.5 kcal/mol to 5.0 kcal/mol. The value of $P(\theta_r)$ at $\theta_r = 90^\circ$ becomes higher with the collision energies change from 2.5 kcal/mol to 4.0 kcal/mol and the highest appears at 4.0 kcal/mol while the value becomes lower with the collision energies changing from 4.0 kcal/mol to 5.0 kcal/mol, which means the change of the product alignment is not a monotonic enhancement with increasing collision energies for $\text{Li+TF} \rightarrow \text{LiF+T}$ reaction. Moreover, there is a prominent feature observed in Figure 3.1 (a) and (b) that the $P(\theta_r)$ value of the $\text{Li+DF} \rightarrow \text{LiF+D}$ reaction at $\theta_r = 90^\circ$ is larger than that of the $\text{Li+TF} \rightarrow \text{LiF+T}$ reaction at the same collision energy, which means that the product alignment of the $\text{Li+DF} \rightarrow \text{LiF+D}$ reaction is stronger than the $\text{Li+TF} \rightarrow \text{LiF+T}$ reaction. As discussed in Refs [27, 29] the $P(\theta_r)$ is sensitive to two factors: one is the character of the PES and the other is the mass factor $\cos^2\beta$ (i.e. $\cos^2\beta = m_A m_C / (m_A + m_B)(m_B + m_C)$) for the $\text{A+BC} \rightarrow \text{AB+C}$ reaction. For reactions $\text{Li+DF/TF} \rightarrow \text{LiF+D/T}$, the mass

factor $\cos^2\beta$ is 0.02585 and 0.03695 respectively. Since the same PES is adopted in the calculation, the change of the $P(\theta_r)$ distributions is due to the difference in mass factor of the reactions, that is, the $P(\theta_r)$ distributions are affected by isotope effect. The product rotational angular momentum is approximately equal to the reactant orbital angular momentum, and the reactant orbital angular momentum plays an important role on the distribution of $P(\theta_r)$ for HHL mass combination. [29] Therefore, large mass factor of the reactions will have larger product rotational angular momentum, leading to the fact that the product rotational angular momentum alignment effect weakens with the mass factor increases. The dihedral angle of $P(\phi_r)$ distributions, which describes the $k - k' - j'$ correlation, shown in Figure 3.1 (c) and (d), can provide both product alignment and orientation information. From Figure 3.1 (c) and (d), it is clear that $P(\phi_r)$ tends to be asymmetric with respect to the scattering plane, reflecting the strong polarization of angular momentum for $\text{Li+DF} \rightarrow \text{LiF+D}$ and $\text{Li+TF} \rightarrow \text{LiF+T}$ reactions at different collision energies. At the same time, the peaks of the product $P(\phi_r)$ distributions of $\text{Li+DF} \rightarrow \text{LiF+D}$ reaction at $\phi_r = 270^\circ$ become higher with the collision energy increase from 2.5 kcal/mol to 5.0 kcal/mol. However, for $\text{Li+TF} \rightarrow \text{LiF+T}$ reaction, while the value of the $P(\phi_r)$ distribution at $\phi_r = 270^\circ$ becomes higher when the collision energies change from 2.5 kcal/mol to 4.5 kcal/mol, the peak of $P(\phi_r)$ distribution at $\phi_r = 270^\circ$ at 5.0 kcal/mol is lower than that of 4.5 kcal/mol, which means the peaks of the $P(\phi_r)$ distributions for $\text{Li+TF} \rightarrow \text{LiF+T}$ reaction are not monotonic enhancement over the collision energies range of 2.5 kcal/mol to 5.0 kcal/mol. Furthermore, the values of all reactions at $\phi_r = 270^\circ$ are much larger than at $\phi_r = 90^\circ$, which reflects that the product rotational angular momentum is oriented along the negative y-axis and the product has a preference for left-handed rotation in planes which are parallel to the scattering plane. Generally speaking, as the collision energy increases, the values of the product $P(\phi_r)$ distribution of $\text{Li+DF/TF} \rightarrow \text{LiF+D/T}$ reactions become higher. Similar to $P(\theta_r)$ distributions, the peaks of $P(\phi_r)$ distributions for $\text{Li+DF} \rightarrow \text{LiF+D}$ reaction are larger than that for $\text{Li+TF} \rightarrow \text{LiF+T}$ reaction at the same collision energy, which indicates that the degree of product orientation of the $\text{Li+DF} \rightarrow \text{LiF+D}$ reaction is stronger than that of $\text{Li+TF} \rightarrow \text{LiF+T}$ reaction. In this case, as discussed in Refs.[32, 33] we can deduce that the discrepancy of the orientation degree for the reaction may due to the different effective potential well depth of the reaction induced by the different harmonic zero point energy (ZPE) of the reactant molecules DF/TF. In the light of $\nu = \sqrt{k/\mu}/2\pi$ and $E_0 = h\nu/2$, the vibrational frequency of DF is higher than that of TF. Consequently, the depth of effective potential well for the $\text{Li+DF} \rightarrow \text{LiF+D}$ reaction is shallower than that of $\text{Li+TF} \rightarrow \text{LiF+T}$ reaction, which makes the degree of product orientation for the $\text{Li+DF} \rightarrow \text{LiF+D}$ reaction stronger. Such phenomenon can also be well explained qualitatively by the impulse model for the $\text{A+BC} \rightarrow \text{AB+C}$ reaction.[33-35] According to Refs.[33-35] we have $j' = L \sin^2\beta + j \cos^2\beta + J_1 m_B / m_{AB}$, where L is the reagent orbital angular momentum and $J_1 = \sqrt{\mu_{BC} R} (r_{AB} + r_{CB})$, with r_{AB} and r_{CB} being the unit vectors for B pointing to A and C , respectively, μ_{BC} is the reduced mass of the BC molecule and R is the repulsive energy between B and C atoms. Moreover, $\cos^2\beta$ is mass factor. During the whole reactive process for the $\text{Li+DF/TF} \rightarrow \text{LiF+D/T}$ reactions, since the term $L \sin^2\beta + j \cos^2\beta$ is symmetric, the term