

ARTIFICIAL BOUNDARY CONDITIONS FOR “VORTEX IN CELL” METHOD^{*1)}

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Dedicated to the 80th birthday of Professor Zhou Yulin

Abstract

This paper mainly designs artificial boundary conditions for “vortex in cell” method in solving two-dimensional incompressible inviscid fluid under two conditions: one is with periodical initial value in one direction and the other with compact supported initial value. To mimic the vortex motion, Euler equation is transformed into vorticity-stream function and the technique of vortex in cell is applied incorporating with the artificial boundary conditions.

Key Words: Incompressible inviscid flow, Vortex in cell method, Artificial boundary condition

1. Mathematical Model

The motion of two-dimensional incompressible inviscid fluid satisfies the following equations:

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} + \frac{1}{\rho} \nabla P = \vec{f} \quad (1.1)$$

$$\nabla \cdot \mathbf{u} = 0, \quad (1.2)$$

where $\mathbf{u}(u, v)$ is the velocity vector of a particle at $\vec{x} = (x, y)$ and time t . $P(\vec{x}, t)$ and $\rho(\vec{x}, t)$ denote the pressure and the density of the fluid respectively. $\rho(x, t)$ is constant for incompressible fluid. Fone \vec{f} acts on a unit fluid. ∇ is gradient operator. ‘ \cdot ’ stands for the inner product. In (x, y) -plane $\nabla \Lambda = (\frac{\partial}{\partial y}, -\frac{\partial}{\partial x})$ is rotation operator. Here we assume $\vec{f} = \nabla \varphi$, where φ is a scalar function, and define vorticity ω as $\omega = -\nabla \Lambda \mathbf{u} = -\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$. Inserting $\nabla \Lambda$ into Euler equation (1.1), one can obtain:

$$\frac{\partial \omega}{\partial t} + (\mathbf{u} \cdot \nabla) \omega = 0. \quad (1.3)$$

Let \vec{x}_0 be an arbitrary point in the plane and define the stream function $\psi(\vec{x})$ as

$$\psi(\vec{x}) = \int_{\vec{x}_0}^{\vec{x}} v dx + u dy,$$

where the integration in (x, y) plane is independent of the integration path due to (??) and Green formula. Applying $\nabla \Lambda$ and $-\nabla \Lambda$ to $\psi(\vec{x})$ in succession, one can get the relationship of

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the stream function and vorticity, i. e., $-\Delta\psi = \omega$. Thus the incompressible inviscid fluid in vorticity-stream function is:

$$\frac{\partial\omega}{\partial t} + u\frac{\partial\omega}{\partial x} + v\frac{\partial\omega}{\partial y} = 0 \quad (1.4)$$

$$-\Delta\psi = \omega \quad (1.5)$$

$$\frac{\partial\psi}{\partial y} = u \quad (1.6)$$

$$\frac{\partial\psi}{\partial x} = -v. \quad (1.7)$$

2. Technique of “vortex in cell”

Christansen [?] is the first one to use the technique of vortex in cell to simulate the motion of two-dimensional incompressible inviscid fluid. This paper adopts this method to simulate vortices interaction in (x, y) -plane. We discrete the fluid vorticity ω into sum of vortex points:

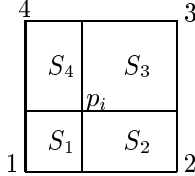
$$\omega(x, y) \approx \sum_{i=1}^M p_i \delta(x - x_i) \delta(y - y_i),$$

where M is the total number of vortex points, p_i is vorticity of the i -th point at (x_i, y_i) . The velocity $\mathbf{u}_i = (u_i, v_i)$ of the i -th vortex point is given by

$$\frac{\partial x_i}{\partial t} = u_i \quad (2.1)$$

$$\frac{\partial y_i}{\partial t} = v_i. \quad (2.2)$$

Once we know the velocity (u_i, v_i) of the vortex point, then we can advance the vortex point by using (2.1)-(2.2). Vortex in cell method solves the Poisson’s equation (1.5) on a uniform mesh by the usual five-point scheme. Thus to mimic the motion of vortex points, it is necessary to allocate the vortex point’s vorticity to mesh points first and then, after solving the Poisson’s equation (1.5) on uniform mesh, to redistribute mesh point’s velocity to vortex points. Let s be the area of an mesh and s_1, s_2, s_3, s_4 be the four parts divided by a vortex point located in the mesh showed by the following figure. Then the vortex point allocates its vorticity p_i to the four surrounding mesh points as follows:



$$\omega_1 = \frac{s_3}{s} p_i \quad \omega_2 = \frac{s_4}{s} p_i \quad \omega_3 = \frac{s_1}{s} p_i \quad \omega_4 = \frac{s_2}{s} p_i.$$

All other vortex points located in the same mesh allocate their vorticity to the mesh points similarly. On the other hand, mesh points distribute their velocity to the vortex point p_i by the technique of vortex in cell. That is, assume \mathbf{v}_l , $l = 1, 2, 3, 4$ are the velocities of the four surrounding points of a mesh, then the velocity \mathbf{u}_i of the vortex point p_i located in the mesh is defined by

$$\mathbf{u}_i = \frac{s_3}{s} \mathbf{v}_1 + \frac{s_4}{s} \mathbf{v}_2 + \frac{s_1}{s} \mathbf{v}_3 + \frac{s_2}{s} \mathbf{v}_4.$$