

# NUMERICAL SOLUTION OF RADON'S PROBLEM IN A TWO DIMENSIONAL SPACE\*

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## § 1

As in the Fourier transform of a function, we associate with a function  $f(x)$  its Radon transform  $g(\alpha, p)$ , defined by the following integral of  $f(x)$  over the hyperplane with unit normal  $\alpha$  and distance  $p$  from the origin:

$$Rf = \int_{x \cdot \alpha = p} f(x) d\omega_\alpha = g(\alpha, p) \quad (1)$$

for  $p \in R^1$ ,  $x, \alpha \in R^n$ ,  $|\alpha| = 1$ . The Radon problem consists in solving equation (1) for  $f(x)$  from  $g(\alpha, p)$ . This problem is of great importance in many applications, for instance, in the reconstruction of objects from X-ray pictures ([1], [2]).

In this paper we shall merely treat Radon's problem for  $n=2$ . Describing the unit normal  $\alpha$  by its polar angle  $\theta$ , we can rewrite (1) as

$$Rf = \int_{-\infty}^{\infty} f(\theta; p, r) dr = g(\theta, p), \quad (2)$$

$$f(\theta; p, r) = f(p \cos \theta + r \sin \theta, p \sin \theta - r \cos \theta),$$

or

$$4\pi \int_p^\infty \frac{\eta a(\eta)}{\sqrt{\eta^2 - p^2}} d\eta = G(p), \quad G(p) = \int_0^{2\pi} g(\theta, p) d\theta,$$

where  $a(\eta)$  is the average of  $f$  on the circle of radius  $\eta$  about the origin:

$$a(\eta) = \frac{1}{\omega \eta^{n-1}} \int_{|\alpha|=1} f(x) ds_\alpha.$$

The problem of determining the solution  $a(\eta)$  (in particular,  $f(\eta)$ , if the function  $f$  has the property of circular symmetry<sup>[1]</sup>, i.e.  $f(x_1, x_2) = f(\eta)$ ,  $x_1^2 + x_2^2 = \eta^2$ ) from the initial data  $G(p)$  has been explored in [3] in greater detail.

Radon's problem (2) is not well-posed on the pair of spaces  $(\bar{C}, L_2)$ <sup>[4]</sup>, where

$$L_2 = L_2(H),$$

$$\begin{aligned} \bar{C} = \bar{C}(K_T) &= \{f(x): f(x) \text{ is continuous and has compact support } K_\xi, 0 < \xi \leq T\}, \\ H &= \{(\alpha, p): p \in R^1, \alpha \in R^2, |\alpha| = 1\}. \end{aligned}$$

is the unit cylinder in  $R^3$  and  $K_\xi$  is the circle of radius  $\xi$  about the origin. This is because the range of Radon's integral operator  $R$  clearly does not coincide with  $L_2$  and the inverse  $R^{-1}$  of the operator  $R$  is not continuous.

It should be pointed out that the reciprocity formula for  $f(x)$  holds ([1], [2]):

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$$(2\pi i)^n f(z) = (4s)^{\frac{n-2}{2}} \int_{D_s} d\omega_s \int_{p=-\infty}^{p=\infty} \frac{dg(x, p)}{p - z \cdot x}.$$

## § 2

Below, we shall use the finite-difference method to solve the Radon equation (2). In order to clarify the essentials of the method, in this section we study merely a semi-discrete scheme, where only the variable  $p$  is discretized:

$$p_i = i\tau, \quad i = -n-1, -n, \dots, n, n+1, \quad \tau = \frac{T}{n+1}$$

while the variable  $r$  is left continuous. The fully discrete case is discussed in section 3.

Let  $f_\tau$  be a vector-valued function:

$$f_\tau = f_\tau(\theta, r) = (f_{-n-1}(\theta, r), \dots, f_{n+1}(\theta, r))$$

the components of which, i.e.,  $f_i(\theta, r)$ , are defined on the segments:

$$\{(p_i, r) : -T \leq r \leq T\} \quad (i = -n-1, \dots, n+1).$$

Now let us consider the functional  $M_\tau^\alpha$  of the argument function  $f_\tau$ , which may be taken to be an arbitrary continuous function with a continuous derivative:

$$M_\tau^\alpha[\theta; f_\tau, g_\tau] = \sum_{i=-n}^n \tau \left[ \int_{-\bar{T}}^{\bar{T}} f_i(\theta, r) dr - g_i(\theta) \right]^2 + \alpha \sum_{i=-n}^n \tau \int_{-\bar{T}}^{\bar{T}} \left\{ f_i^2(\theta, r) + \left[ \frac{df_i}{dr} \right]^2 \right\} dr \\ + \alpha \sum_{i=-n-1}^n \tau \int_{-\bar{T}}^{\bar{T}} \left[ \frac{f_{i+1} - f_i}{\tau} \right]^2 dr,$$

where  $g_\tau$  is a given vector:

$$g_\tau = g_\tau(\theta) = (g_{-n}(\theta), \dots, g_n(\theta)).$$

**Theorem 1.** For every  $g_\tau$  and every positive parameter  $\alpha$ , there exists a unique continuous function  $f_\tau^\alpha(\theta, r)$  with a continuous derivative for which the functional  $M_\tau^\alpha[\theta; f_\tau, g_\tau]$  attains its greatest lower bound:

$$M_\tau^\alpha[\theta; f_\tau^\alpha, g_\tau] = \inf M_\tau^\alpha[\theta; f_\tau, g_\tau].$$

*Proof.* 1) The desired function  $f_\tau^\alpha(\theta, r)$  should be determined by the integro-differential equation of Euler:

$$\alpha L[f_\tau] = \int_{-\bar{T}}^{\bar{T}} f_\tau(\theta, s) ds - g_\tau(\theta), \quad (3)$$

$$L[f_\tau] = \frac{d^2 f_\tau}{dr^2} + B f_\tau,$$

$$B = \begin{pmatrix} -\left(1 + \frac{1}{\tau^2}\right) & \frac{1}{\tau^2} \\ \frac{1}{\tau^2} & -\left(1 + \frac{1}{\tau^2}\right) \end{pmatrix}$$