A Study on Oscillatory and Asymptotic Nature of Impulsive Neutral Differential Equations of Order Three^{*}

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Abstract In this paper, we consider a class of third-order neutral impulsive differential equations. An equivalent class of neutral differential equations is obtained by using a suitable substitution. Some new oscillation results are proved. Moreover, we discuss the asymptotic behavior of the solution. The results presented here are illustrated via examples.

Keywords Neutral differential equations, Impulsive conditions, Oscillation criteria, Asymptotic behavior.

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1. Introduction

Here, we establish the oscillation results for the following model of neutral impulsive differential equation of order three:

$$\begin{cases} \left(s_1(t)\upsilon^{(2)}(t)\right)^{(1)} + s_2(t)u^{(1)}(\eta_2(t)) = 0, \quad t \neq t_p, \\ u^{(r)}(t_p) - u^{(r)}(t_p^-) = d_p u^{(r)}(t_p^-), \quad r = 0, 1, 2, \\ p = 1, 2, 3, \dots, \end{cases}$$
(1.1)

where $v(t) = u(t) + \alpha u(\eta_1(t)), \eta_1(t) \leq t, \eta_2(t) \leq t, t > t_0, \alpha > 0$, and $v^{(r)}(t)$ denotes the derivative of order r with respect to t.

It is well-known that the motions on the earth are not always uniform, as various kinds of resistance appear during the motions. For example, suppose high-intensity forces act for a short duration of time. In that case, motions caused by these forces are called impulsive motions, and the differential equations describing these motions are called impulsive differential equations.

The differential equations with impulsive effect can be used to simulate those discontinuous processes in which impulses occur. Therefore, it has become an important tool to handle the real process of mathematical models and phenomena

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such as optimal control, electric circuit, biotechnology, population dynamics, fractals, neural network, viscoelasticity and chemical technology. For more details on impulsive differential equations, please refer to [15]. One of the main advantages of the impulses can be seen in [23], as they provided the model in which the mass point might oscillate in the presence of impulsive effect and in the absence of impulsive effect, the mass point did not oscillate. For more work on impulsive effect, please refer to [7,14,19].

In 1989, some researchers started to investigate the oscillatory nature of differential equations with impulses, and they were at the initial stage of development. Later on, the authors of [7,9-11,18,23] extended the study of oscillation to parabolic and hyperbolic impulsive partial differential equations. A hybrid evolution system with impulsive conditions has been studied by Sadhasivam and Deepa [21]. In the last few decades, many researchers [2, 4, 13, 20, 22, 25, 29] have applied the Riccati technique to the study of the oscillatory behavior for various types of second-order differential equations. Some works on the oscillatory and asymptotic behavior of the solutions to higher-order impulsive differential equations have been carried out in [8, 16, 27]. Li [16] investigated the oscillatory and asymptotic behavior of the solutions to a higher-order delay differential equation with impulses by using comparison results with an associated non-impulsive delay differential equation.

Basic definitions and results on oscillation for neutral type differential equations can be found in [5,15]. Due to the wide applicability of neutral differential equations in various fields of science and engineering, there is a great interest in obtaining new oscillation criteria for a different types of differential equations (see, for instance, [2–4,6,12,13,17,22,25,26,29]). We have often seen that even non-impulsive neutral delay differential equations may have solutions of oscillatory nature due to some additional control. An improved sufficient condition for the oscillation and asymptotic stability has been obtained in this paper [26]. Guan and Shen [12] examined the oscillation criteria of a first-order impulsive differential equations with variable delays. Oscillation theorems for third-order delay differential equations were discussed by Tiryaki and Aktas [24]. Arul and Shobha [2] generalized neutral differential equations of order two and presented some new oscillation criteria by using Riccati transformation under some conditions.

Recently, Zhang and Li [28] have studied the oscillatory behaviour of the solutions to second-order impulsive neutral dynamic system with positive and negative coefficients. Moreover, in [22], the authors have presented some new necessary and sufficient conditions for the oscillation of a class of second-order neutral delay impulsive differential equations.

Motivated by all the above works, we obtain some new oscillation results of impulsive neutral differential equations of order three by converting them to the non-impulsive neutral differential equations. Also, there are only a few papers that deal with this technique.

The rest of this paper is organized as follows. Section 2 contains some basic lemmas and assumptions, which are required for the next sections. In Section 3, some new-type oscillation results are obtained for problem (1.1) by using Riccati transformation. In the last section, the results are illustrated by examples.

2. Preliminaries and assumptions

Throughout the paper, we consider the following assumptions:

- (C1) The functions $s_r: (t_0, \infty) \to \mathbb{R}^+, r = 1, 2$ are continuous, and there exists G > 0 such that $s_1(t) \leq G$.
- (C2) $\eta_r: (t_0, \infty) \to \mathbb{R}, r = 1, 2$ are continuous functions with the following conditions:
 - (i) $\eta_1(t) \le t, \, \eta_2(t) \le t, \, \eta_1(\eta_2(t)) = \eta_2(\eta_1(t));$

 - (ii) $\eta_1^{(1)}(t) = \eta_2^{(1)}(t) = 1;$ (iii) $\lim_{t \to \infty} \eta_1(t) = \infty, \lim_{t \to \infty} \eta_2(t) = \infty.$

Lemma 2.1. $\zeta(t) = \prod_{t_0 < t_p \le t} (1 + d_p)^{-1} u(t)$ satisfies

$$\left(s_1(t)V^{(2)}(t)\right)^{(1)} + s_2(t)\prod_{\eta_2(t) < t_p \le t} (1+d_p)^{-1}\zeta^{(1)}(\eta_2(t)) = 0,$$
(2.1)

where

$$V(t) = \zeta(t) + \alpha \prod_{\eta_1(t) < t_p \le t} (1 + d_p)^{-1} \zeta(\eta_1(t)), \qquad (2.2)$$

if u(t) satisfies (1.1) on the interval (t_0, ∞) .

Proof. Let $\zeta(t) = \prod_{t_0 < t_p \le t} (1 + d_p)^{-1} u(t)$ satisfy (2.1). Then, we will show that u(t) satisfies (1.1) on the interval (t_0, ∞) . Obviously,

$$u(t) = \prod_{t_0 < t_p \le t} (1+d_p)\zeta(t) \text{ and } v(t) = \prod_{t_0 < t_p \le t} (1+d_p)V(t).$$

Thus,

$$\left(s_1(t)v^{(2)}(t)\right)^{(1)} = \prod_{t_0 < t_p \le t} (1+d_p) \left(s_1(t)V^{(2)}(t)\right)^{(1)}.$$

Using (2.1), we get

$$\left(s_1(t) \upsilon^{(2)}(t) \right)^{(1)} = \prod_{t_0 < t_p \le t} (1+d_p) \left(-s_2(t) \prod_{\eta_2(t) < t_p \le t} (1+d_p)^{-1} \zeta^{(1)}(\eta_2(t)) \right)$$

= $-s_2(t) \prod_{t_0 < t_p \le \eta_2(t)} (1+d_p) \zeta^{(1)}(\eta_2(t)).$

Therefore, for $t \neq t_p$, we have

$$\left(s_1(t)v^{(2)}(t)\right)^{(1)} + s_2(t)u^{(1)}(\eta_2(t)) = 0.$$

Also, we obtain

$$u^{(r)}(t) = \prod_{t_0 < t_p \le t} (1+d_p)\zeta^{(r)}(t), \quad r = 0, 1, 2.$$

$$\implies u^{(r)}(t_p) = (1+d_p)u^{(r)}(t_p^-).$$

This shows that u(t) satisfies (1.1).

Conversely, assume that $u(t) = \prod_{t_0 < t_p \le t} (1 + d_p)\zeta(t)$ satisfies (1.1). Then, we will show that $\zeta(t)$ satisfies (2.1) on (t_0, ∞) . As $V(t) = \prod_{t_0 < t_p \le t} (1 + d_p)^{-1} \upsilon(t)$, we have

$$\left(s_1(t)V^{(2)}(t)\right)^{(1)} = \prod_{t_0 < t_p \le t} (1+d_p)^{-1} \left(s_1(t)v^{(2)}(t)\right)^{(1)}.$$

Using (1.1), we obtain

$$\left(s_1(t) V^{(2)}(t) \right)^{(1)} = -s_2(t) \prod_{\substack{t_0 < t_p \le t}} (1+d_p)^{-1} u^{(1)}(\eta_2(t))$$
$$= -s_2(t) \prod_{\substack{\eta_2(t) < t_p \le t}} (1+d_p)^{-1} \zeta^{(1)}(\eta_2(t)).$$

On the other hand, we can easily show that $\zeta^{(r)}(t_p^-) = \zeta^{(r)}(t_p)$. This shows that $\zeta(t)$ satisfies (2.1).

Lemma 2.2 ([16]). A non-zero solution u(t) of (1.1) is oscillatory on (t_0, ∞) , if the corresponding solution $\zeta(t) = \prod_{t_0 < t_p \le t} (1 + d_p)^{-1} u(t)$ of (2.1) is oscillatory on (t_0, ∞) . Moreover, $\lim_{t \to \infty} u^{(r)}(t) = 0$, if $\lim_{t \to \infty} \zeta^{(r)}(t) = 0$, r = 0, 1, 2.

3. Main results

Theorem 3.1. Suppose one of the following three conditions:

$$\int_{t_0}^{\infty} \nu^2 s_2(\nu) \Upsilon(\nu) d\nu = -\infty, \qquad (3.1)$$

$$\int_{t_0}^t J(\nu)\Upsilon(\nu)d\nu = \infty$$
(3.2)

and

$$\int_{t_0}^t \left(S_1(\nu) J(\nu) \Upsilon(\nu) - \frac{s_1(\nu) + \alpha s_1(\eta_1(\nu))}{4S_1(\nu) s_1^2(\nu)} \right) d\nu = \infty,$$
(3.3)

where

$$J(t) = \min\{s_2(t), s_2(\eta_1(t))\}, \ \Upsilon(t) = \prod_{\eta_2(t) < t_p \le t} (1+d_p)^{-1}, \ S_1(t) = \int_t^\infty \frac{d\nu}{s_1(\nu)}$$

holds. Then, all solutions u(t) of (1.1) such that $u(t)u^{(1)}(t) > 0$ or $u(t)u^{(1)}(t) < 0$, for $t \ge \varrho$ are oscillatory.

Proof. On contrary, assume that $\zeta(t) \neq 0$ be a non-oscillatory solution of (2.1). Further, if we assume that $\zeta(t) > 0$ for $t \ge \rho$, then the following three cases arise.

Case 1: If there exists $\rho \ge t_0$ such that $\zeta(t) > 0$, $\zeta^{(1)}(t) < 0$, $\zeta^{(2)}(t) > 0$, or V(t) > 0, $V^{(1)}(t) < 0$, $V^{(2)}(t) > 0$, for $t \ge \rho$. Since $\zeta^{(2)}(\eta_2(t)) > 0$, $\zeta^{(1)}(\eta_2(t))$ is increasing for $t \ge \rho$. Therefore, for $t \ge \rho$, from (2.1), we get

$$\left(s_1(t) V^{(2)}(t) \right)^{(1)} = -s_2(t) \prod_{\eta_2(t) < t_p \le t} (1+d_p)^{-1} \zeta^{(1)}(\eta_2(t))$$

$$\le -s_2(t) \prod_{\eta_2(t) < t_p \le t} (1+d_p)^{-1} \zeta^{(1)}(\eta_2(\varrho))$$

$$\le -Ms_2(t) \Upsilon(t),$$

where $M = \zeta^{(1)}(\eta_2(\varrho)) < 0$. Integrating ϱ and t after multiplying by ν^2 , we get

$$\int_{\varrho}^{t} \nu^{2} \Big(s_{1}(\nu) V^{(2)}(\nu) \Big)^{(1)} d\nu \leq -M \int_{\varrho}^{t} \nu^{2} s_{2}(\nu) \Upsilon(\nu) d\nu.$$
(3.4)

On the other hand,

$$\begin{split} \int_{\varrho}^{t} \nu^{2} \Big(s_{1}(\nu) V^{(2)}(\nu) \Big)^{(1)} d\nu &= \int_{\varrho}^{t} \nu^{2} d \big(s_{1}(\nu) V^{(2)}(\nu) \big) \\ &= \Big[\nu^{2} s_{1}(\nu) V^{(2)}(\nu) \Big]_{\varrho}^{t} - \int_{\varrho}^{t} 2\nu s_{1}(\nu) V^{(2)}(\nu) d\nu \\ &\geq \Big[\nu^{2} s_{1}(\nu) V^{(2)}(\nu) \Big]_{\varrho}^{t} - 2G \int_{\varrho}^{t} \nu V^{(2)}(\nu) d\nu \\ &= \Big[\nu^{2} s_{1}(\nu) V^{(2)}(\nu) \Big]_{\varrho}^{t} - 2G \Big[\nu V^{(1)}(\nu) \Big]_{\varrho}^{t} \\ &+ 2G \int_{\varrho}^{t} V^{(1)}(\nu) d\nu \\ &= t^{2} s_{1}(t) V^{(2)}(t) - \varrho^{2} s_{1}(\varrho) V^{(2)}(\varrho) \\ &- 2G \big[t V^{(1)}(t) - \varrho V^{(1)}(\varrho) \big] + 2G \big[V(t) - V(\varrho) \big] \\ &\geq -\varrho^{2} s_{1}(\varrho) V^{(2)}(\varrho) + 2G \varrho V^{(1)}(\varrho) - 2G V(\varrho). \end{split}$$

From (3.4), we have

$$-\int_{\varrho}^{t}\nu^{2}s_{2}(\nu)\Upsilon(\nu)d\nu \leq \frac{1}{M}\left\{\varrho^{2}s_{1}(\varrho)V^{(2)}(\varrho) - 2G\varrho V^{(1)}(\varrho) + 2GV(\varrho)\right\} < \infty.$$

On letting $t \to \infty$, we get a contradiction to condition (3.1). Case 2: If there exists $\varrho' > t_0$ such that $\zeta(t) > 0$, $\zeta^{(1)}(t) > 0$, $\zeta^{(2)}(t) > 0$, or V(t) > 0, $V^{(1)}(t) > 0$, $V^{(2)}(t) > 0$, for $t \ge \varrho'$. From (2.1), we have

$$\left(s_1(t) V^{(2)}(t) \right)^{(1)} + \alpha \left(s_1(\eta_1(t)) V^{(2)}(\eta_1(t)) \right)^{(1)} + J(t) \Upsilon(t) \\ \times \left[\zeta^{(1)}(\eta_2(t)) + \alpha \prod_{\eta_1(\eta_2(t)) < t_p \le \eta_2(t)} (1+d_p)^{-1} \zeta^{(1)}(\eta_1(\eta_2(t))) \right] \le 0.$$

Using (2.2), for $t \ge \varrho'$, we obtain

$$\left(s_1(t)V^{(2)}(t)\right)^{(1)} + \alpha \left(s_1(\eta_1(t))V^{(2)}(\eta_1(t))\right)^{(1)} + J(t)\Upsilon(t)V^{(1)}(\eta_2(t)) \le 0.$$
(3.5)

Using the fact $V^{(1)}(\eta_2(t)) \ge c > 0$ and integrating from ϱ' to t, we get

$$\int_{\varrho'}^t J(\nu) \Upsilon(\nu) d\nu < \infty,$$

which is a contradiction to condition (3.2).

Case 3: If there exists $\rho'' \ge t_0$ such that $\zeta(t) > 0$, $\zeta^{(1)}(t) > 0$, $\zeta^{(2)}(t) < 0$, or V(t) > 0, $V^{(1)}(t) > 0$, $V^{(2)}(t) < 0$, for $t \ge \rho''$. Define

$$\chi(t) = \frac{s_1(t)V^{(2)}(t)}{V^{(1)}(t)}, \quad t \ge \varrho''.$$
(3.6)

Since $s_1(\nu)V^{(2)}(\nu)$ is decreasing, for $\nu \ge t$, we have

$$s_1(\nu)V^{(2)}(\nu) \le s_1(t)V^{(2)}(t).$$

$$\implies V^{(2)}(\nu) \le s_1(t)\frac{V^{(2)}(t)}{s_1(\nu)}$$

Integrating from t to μ , we get

$$V^{(1)}(\mu) - V^{(1)}(t) \le s_1(t)V^{(2)}(t) \int_t^{\mu} \frac{d\nu}{s_1(\nu)}, \quad \mu \ge t \ge \varrho''.$$

Taking limit as $\mu \to \infty$, we get

$$\frac{s_1(t)V^{(2)}(t)}{V^{(1)}(t)}S_1(t) > -1, \quad t \ge \varrho''.$$

$$\implies -1 < \chi(t)S_1(t) < 0, \quad t \ge \varrho''.$$
(3.7)

On the other hand, we define

$$\kappa(t) = \frac{s_1(\eta_1(t))V^{(2)}(\eta_1(t))}{V^{(1)}(t)}, \quad t \ge \varrho''.$$
(3.8)

Since $s_1(t)V^{(2)}(t)$ is decreasing, we have $s_1(t)V^{(2)}(t) \le s_1(\eta_1(t))V^{(2)}(\eta_1(t))$. Then, $\chi(t) \le \kappa(t)$. From (3.7), we have

$$-1 < \kappa(t)S_1(t) < 0, \quad t \ge \varrho''.$$
 (3.9)

Differentiating (3.6), we get

$$\chi^{(1)}(t) = \frac{\left(s_1(t)V^{(2)}(t)\right)^{(1)}}{V^{(1)}(t)} - \frac{s_1(t)V^{(2)}(t)}{(V^{(1)}(t))^2}V^{(2)}(t)$$
$$= \frac{\left(s_1(t)V^{(2)}(t)\right)^{(1)}}{V^{(1)}(t)} - \frac{\chi^2(t)}{s_1(t)}, \quad t \ge \varrho''. \tag{3.10}$$

Differentiating (3.8), we get

$$\kappa^{(1)}(t) \le \frac{\left(s_1(\eta_1(t))V^{(2)}(\eta_1(t))\right)^{(1)}}{V^{(1)}(t)} - \frac{Y^2(t)}{s_1(\eta_1(t))}, \quad t \ge \varrho''.$$
(3.11)

Using (3.10) and (3.11), we get

$$\chi^{(1)}(t) + \alpha \kappa^{(1)}(t) \leq \frac{\left(s_1(t)V^{(2)}(t)\right)^{(1)}}{V^{(1)}(t)} + \alpha \frac{\left(s_1(\eta_1(t))V^{(2)}(\eta_1(t))\right)^{(1)}}{V^{(1)}(t)} - \frac{\chi^2(t)}{s_1(t)} - \alpha \frac{\kappa^2(t)}{s_1(\eta_1(t))}.$$

Using (3.5), we get

$$\chi^{(1)}(t) + \alpha \kappa^{(1)}(t) \le -J(t)\Upsilon(t) - \frac{\chi^2(t)}{s_1(t)} - \alpha \frac{\kappa^2(t)}{s_1(\eta_1(t))}.$$

Integrating the above inequality from ρ'' to t after multiplying by $S_1(t)$, we get

$$\int_{\varrho''}^t S_1(\nu) \Big[\chi^{(1)}(\nu) + \alpha \kappa^{(1)}(\nu) \Big] d\nu \leq -\int_{\varrho''}^t S_1(\nu) J(\nu) \Upsilon(\nu) d\nu -\int_{\varrho''}^t S_1(\nu) \left(\frac{\chi^2(\nu)}{s_1(\nu)} + \alpha \frac{\kappa^2(\nu)}{s_1(\eta_1(\nu))} \right) d\nu.$$

After some simplification, by using (3.7) and (3.9), we obtain

$$-(1+\alpha) + \int_{\varrho''}^t \left(\frac{S_1(\nu)\chi^2(\nu)}{s_1(\nu)} + \frac{\chi(\nu)}{s_1(\nu)}\right) d\nu - S_1(\varrho'') \left(\chi(\varrho'') + \alpha\kappa(\varrho'')\right)$$
$$+\alpha \int_{\varrho''}^t \left(\frac{S_1(\nu)\kappa^2(\nu)}{s_1(\eta_1(\nu))} + \frac{\kappa(\nu)}{s_1(\nu)}\right) d\nu < -\int_{\varrho''}^t S_1(\nu)J(\nu)\Upsilon(\nu)d\nu.$$

Using the inequality $\lambda AB^{\lambda-1} - A^{\lambda} \leq (\lambda - 1)B^{\lambda}, \ \lambda \geq 1$, we obtain

$$\begin{split} \int_{\varrho''}^t \left(S_1(\nu) J(\nu) \Upsilon(\nu) - \frac{s_1(\nu) + \alpha s_1(\eta_1(\nu))}{4S_1(\nu) s_1^2(\nu)} \right) d\nu \\ &\leq (1+\alpha) + S_1(\varrho'') \Big(\chi(\varrho'') + \alpha \kappa(\varrho'') \Big) < \infty, \end{split}$$

which is a contradiction to condition (3.3). Applying Lemma 2.2, the result follows. $\hfill \Box$

Theorem 3.2. Let u(t) be an eventually positive solution of (1.1) such that $u^{(1)}(t) > 0$, $u^{(2)}(t) < 0$, and there exists a constant K > 0 such that

$$\prod_{t_0 < t_p \le t} (1 + d_p)^{-1} \le K.$$

Further, if

$$\int_{t}^{\infty} \frac{1}{s_{1}(\nu_{1})} \left[\int_{\varrho}^{\nu_{1}} s_{2}(\nu) d\nu \right] d\nu_{1} = \infty, \qquad (3.12)$$

then $\lim_{t \to \infty} u^{(1)}(t) = 0.$

Proof. Since u(t) has the property u(t) > 0, $u^{(1)}(t) > 0$, $u^{(2)}(t) < 0$, or $\zeta(t) > 0$, $\zeta^{(1)}(t) > 0$, $\zeta^{(2)}(t) < 0$, and consequently V(t) > 0, $V^{(1)}(t) > 0$, $V^{(2)}(t) < 0$, for $t \ge \varrho$. Let $\lim_{t\to\infty} V^{(1)}(t) = A \ge 0$. Then, we need to prove that A = 0. On contrary, assume A > 0. Since $V^{(1)}(t)$ is decreasing for $t \ge \varrho$, we have $A + \epsilon > V^{(1)}(t) > A$, for all $\epsilon > 0$. Choosing $0 < \epsilon < \frac{A(1-\alpha K)}{\alpha K}$, it is easy to see that

$$\begin{aligned} \zeta^{(1)}(t) &= V^{(1)}(t) - \alpha \prod_{\eta_1(t) < t_p \le t} (1 + d_p)^{-1} \zeta^{(1)}(\eta_1(t)) \\ &\ge A - \alpha \prod_{\eta_1(t) < t_p \le t} (1 + d_p)^{-1} V^{(1)}(\eta_1(t)) \\ &\ge A - \alpha \prod_{\eta_1(t) < t_p \le t} (1 + d_p)^{-1} (A + \epsilon) \\ &\ge A - \alpha K(A + \epsilon) = P(A + \epsilon) > PV^{(1)}(t), \end{aligned}$$
(3.13)

where $P = \frac{A - \alpha K(A + \epsilon)}{A + \epsilon} > 0.$

Therefore, from (2.1), we have

$$\left(s_1(t) V^{(2)}(t) \right)^{(1)} = -s_2(t) \prod_{\eta_2(t) < t_p \le t} (1+d_p)^{-1} \zeta^{(1)}(\eta_2(t))$$

$$\le -s_2(t) K P V^{(1)}(t)$$

$$\le -s_2(t) K P A.$$

Integrating from ρ to t, we get

$$s_1(t)V^{(2)}(t) \le -KPA \int_{\varrho}^{t} s_2(\nu)d\nu.$$
$$\implies V^{(2)}(t) \le -\frac{KPA}{s_1(t)} \int_{\varrho}^{t} s_2(\nu)d\nu.$$

Again, integrating from t to ∞ , we get

$$-V^{(1)}(t) \leq -KPA \int_t^\infty \frac{1}{s_1(\nu_1)} \left[\int_{\varrho}^{\nu_1} s_2(\nu) d\nu \right] d\nu_1.$$

$$\implies V^{(1)}(t) \geq KPA \int_t^\infty \frac{1}{s_1(\nu_1)} \left[\int_{\varrho}^{\nu_1} s_2(\nu) d\nu \right] d\nu_1.$$

which contradicts condition (3.12). Hence, $\lim_{t \to \infty} V^{(1)}(t) = 0$.

Since $\zeta^{(1)}(t) \leq V^{(1)}(t)$, we have

$$\lim_{t \to \infty} \zeta^{(1)}(t) = 0.$$

Therefore, by applying Lemma 2.2, we get

$$\lim_{t \to \infty} u^{(1)}(t) = 0.$$

4. Applications

In this section, we provide some examples to interpret the main results.

Example 4.1. Consider the given system

$$\begin{cases} \left(t^{-3} \left(u(t) + \alpha u(t-\pi)\right)^{(2)}\right)^{(1)} + t^{-4} u^{(1)}(t-2\pi) = 0, \quad t \neq t_p, \\ u^{(r)}(t_p) - u^{(r)}(t_p^-) = \frac{1}{p} u^{(r)}(t_p^-), \quad r = 0, 1, 2 \quad p = 1, 2, 3, \dots \end{cases}$$
(4.1)

Here, $s_1(t) = t^{-3}$, $s_2(t) = t^{-4}$, $\eta_1(t) = t - \pi$, $\eta_2(t) = t - 2\pi$, $d_p = \frac{1}{p}$, $t_p = p\pi$, $\Upsilon(t) = \prod_{t-2\pi < t_p \le t} \left(\frac{p}{p+1}\right)$. We see that

$$\begin{split} \int_{t_0}^{\infty} \nu^2 s_2(\nu) \Upsilon(\nu) d\nu &= \int_{t_0}^{\infty} \nu^{-2} \prod_{\nu-2\pi < t_p \le \nu} \left(\frac{p}{p+1}\right) d\nu \\ &= \int_{t_0}^{t_1} \nu^{-2} \prod_{\nu-2\pi < t_p \le \nu} \left(\frac{p}{p+1}\right) d\nu + \int_{t_1}^{t_2} \nu^{-2} \prod_{\nu-2\pi < t_p \le \nu} \left(\frac{p}{p+1}\right) d\nu \\ &+ \int_{t_2}^{t_3} \nu^{-2} \prod_{\nu-2\pi < t_p \le \nu} \left(\frac{p}{p+1}\right) d\nu + \cdots \\ &= \frac{1}{2} t_0 - \pi \left[\frac{1}{2} + \frac{1}{3} + \frac{1}{4} + \cdots\right] \\ &= -\infty. \end{split}$$

Thus, condition (3.1) is satisfied. If u(t) > 0, $u^{(1)}(t) < 0$, $u^{(2)}(t) > 0$, for $t \ge \rho$, then by using Case (1) of Theorem 3.1, every non-zero solution u(t) of system (4.1) such that $u(t)u^{(1)}(t) > 0$ or $u(t)u^{(1)}(t) < 0$ is oscillatory.

Example 4.2. Consider the following system

$$\begin{cases} \left(t^{-\frac{1}{2}}\left(u(t) + \alpha u(t-\pi)\right)^{(2)}\right)^{(1)} + tu^{(1)}(t-2\pi) = 0, \quad t \neq t_p, \\ u^{(r)}(t_p) - u^{(r)}(t_p^-) = \frac{1}{p}u^{(r)}(t_p^-), \quad r = 0, 1, 2, \quad p = 1, 2, 3, \dots \end{cases}$$
(4.2)

Here, $s_1(t) = \frac{1}{\sqrt{t}}$, $s_2(t) = t$, $\eta_1(t) = t - \pi$, $\eta_2(t) = t - 2\pi$, $d_p = \frac{1}{p}$, $t_p = p\pi$, $J(t) = \min\{t, t - \pi\} = t - \pi$,

$$\Upsilon(t) = \prod_{\eta_2(t) < t_p \le t} (1+d_p)^{-1} = \prod_{t-2\pi < t_p \le t} \left(\frac{p}{p+1}\right), \quad S_1(t) = \int_t^\infty \sqrt{\nu} d\nu = \infty.$$

We see that

$$\begin{split} \int_{t_0}^{\infty} J(\nu) \Upsilon(\nu) d\nu &= \int_{t_0}^{\infty} (\nu - \pi) \prod_{\nu - 2\pi < t_p \le \nu} \left(\frac{p}{p+1}\right) d\nu \\ &= \frac{1}{2} \int_{t_0}^{t_1} (\nu - \pi) d\nu + \frac{1}{3} \int_{t_1}^{t_2} (\nu - \pi) d\nu + \frac{1}{4} \int_{t_2}^{t_3} (\nu - \pi) d\nu + \cdots \\ &= \frac{1}{2} \left[\frac{(t_1 - \pi)^2}{2} - \frac{(t_0 - \pi)^2}{2} \right] + \frac{1}{3} \left[\frac{(t_2 - \pi)^2}{2} - \frac{(t_1 - \pi)^2}{2} \right] \\ &\quad + \frac{1}{4} \left[\frac{(t_3 - \pi)^2}{2} - \frac{(t_2 - \pi)^2}{2} \right] \\ &= \infty. \end{split}$$

Thus, condition (3.2) is satisfied. If u(t) > 0, $u^{(1)}(t) > 0$, $u^{(2)}(t) > 0$, for $t \ge \varrho$, then by using Case (2) of Theorem 3.1, every non-zero solution u(t) of system (4.1) such that $u(t)u^{(1)}(t) > 0$ or $u(t)u^{(1)}(t) < 0$ is oscillatory. Further, we can easily check that

$$\int_{t_0}^{\infty} \left(S_1(\nu) J(\nu) \Upsilon(\nu) - \frac{s_1(\nu) + \alpha s_1(\eta_1(\nu))}{4S_1(\nu) s_1^2(\nu)} \right) d\nu = \infty.$$

Thus, condition (3.3) is satisfied. If u(t) > 0, $u^{(1)}(t) > 0$, $u^{(2)}(t) < 0$, for $t \ge \varrho$, then by using Case (3) of Theorem 3.1, every non-zero solution u(t) of system (4.1) such that $u(t)u^{(1)}(t) > 0$ or $u(t)u^{(1)}(t) < 0$ is oscillatory.

Example 4.3. Consider the given system

$$\begin{cases} \left(t^{-1}(u(t) + \alpha u(t - \pi))^{(2)}\right)^{(1)} + t^{-4}u^{(1)}(t - 2\pi) = 0, \quad t \neq t_p, \\ u^{(r)}(t_p) - u^{(r)}(t_p^-) = \left(\frac{1+2p}{p^2}\right)u^{(r)}(t_p^-), \quad r = 0, 1, 2 \quad p = 1, 2, 3, \dots \end{cases}$$
(4.3)

Here, $s_1(t) = t^{-1}$, $s_2(t) = t^{-4}$, $\eta_1(t) = t - \pi$, $\eta_2(t) = t - 2\pi$, $d_p = \frac{1+2p}{p^2}$, $t_p = p\pi$. Obviously,

$$\prod_{t_0 < t_p \le t} \left(\frac{p}{p+1}\right)^2 \le K, \quad \text{for some} \quad K > 0.$$

We see that

$$\int_t^\infty \frac{1}{s_1(\nu_1)} \left[\int_{\varrho}^{\nu_1} s_2(\nu) d\nu \right] d\nu_1 = \infty.$$

Thus, condition (3.12) is satisfied. If u(t) > 0, $u^{(1)}(t) > 0$, $u^{(2)}(t) > 0$, for $t \ge \rho$, then by using Theorem 3.2, we get $\lim_{t\to\infty} u^{(1)}(t) = 0$.

5. Conclusion

In this work, we have established some sufficient conditions for the oscillation of solutions for a class of neutral impulsive differential equations of order three, and discussed the asymptotic behavior of solutions. The results are demonstrated by examples.

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