

DECAY OF SOLUTION OF A PARABOLIC EQUATION IN 2-SPACE DIMENSIONS

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Abstract We present a simple method for verifying the uniform L^1 bound and establish sharp rates of L^2 decay of the global solution to the initial value problem for a 2-dimensional parabolic equation.

Key Words Decay estimates; 2D parabolic equation.

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1. Introduction

Newton-Boussinesq equations are of the form

$$\Delta \Psi_t + J(\Delta \Psi, \Psi) = \Delta^2 \Psi - \alpha \theta_x \quad (1)$$

$$\theta_t + J(\theta, \Psi) = \beta \Delta \theta \quad (2)$$

which describe the famous Benard flow. In these last equations, Ψ is flow function, $\Delta \Psi$ is vortex, $\left(\frac{\partial \Psi}{\partial y}, -\frac{\partial \Psi}{\partial x}\right)$ is velocity vector, θ is the temperature, $J(\theta, \Psi) = \frac{\partial \theta}{\partial x} \frac{\partial \Psi}{\partial y} - \frac{\partial \theta}{\partial y} \frac{\partial \Psi}{\partial x}$, $\alpha > 0, \beta > 0$ are constants.

The global existence of the generalized and classical solutions have been established. Spectral method and nonlinear Galerkin methods for solving two-dimensional Newton-Boussinesq equations have been discussed [1-2]. Universal attractors for the Benard problem, existence and physical bounds on their fractal dimension have also been investigated [3].

We note that if we put $\Phi = \Delta \Psi$, and neglect the effect of the temperature in (1), we get

$$\Phi_t + J(\Phi, \Psi) = \Delta \Phi \quad (3)$$

We are interested in the long time behavior of the global solution to the initial value problem for the following parabolic equation

$$\theta_t + J(\theta, \Psi) = \beta \Delta \theta, \quad \beta > 0 \quad (4)$$

$$\theta(x, y, 0) = \theta_0(x, y) \quad (5)$$

where $\theta(x, y, t), \Psi(x, y, t)$ are known scalar functions of the real variables $-\infty < x, y < \infty, 0 \leq t < \infty; \beta > 0$ is a constant.

In this paper, we want to establish sharp L^2 decay of the global solution to problem (1-2), with initial data $\theta_0(x) \in L^1 \cap L^2$. The decay results follow from the *a priori* L^1, L^2 integral estimates and the Fourier transform. The standard argument relies on a technique that involves the splitting of the phase space into two time-dependent domains. For this information, please refer to [4-7].

There has been considerable literature on decay of solutions to the initial value problems for nonlinear evolution equations. Schonbek studied decay of solutions to parabolic conservation laws

$$U_t + \sum_{k=1}^n \frac{\partial}{\partial x_k} f_k(U) = \varepsilon \Delta U, \quad U(x, 0) = U_0(x) \quad (6)$$

She established that if $U_0(x) \in L^1 \cap H^2$, then

$$\int_{R^n} |U|^2 dx \leq C(1+t)^{-n/2}$$

Moreover, under the assumption

$$\left| \frac{d}{dU} f_k(U) \right| \leq C|U|^p, \quad p \geq 1 + \frac{4}{n}, \quad |U| \leq 1$$

or

$$|f_k(U)| \leq C(U)^q, \quad q \geq 2 \left(1 + \frac{1}{n} \right), \quad |U| \leq 1$$

she established the L^∞ optimal decay estimate

$$\|U(t)\|_\infty \leq C(1+t)^{-n/2}$$

Readers who are interested in this problem can find other similar works on solutions to nonlinear evolution equations in our references.

Denote by $Q_t = \{(x, y, s) : -\infty < x, y < \infty, 0 \leq s \leq t\}$, where $0 \leq t < \infty$.

For simplicity, we will denote by C any positive constant which depends only on the norms of the initial data θ_0 , and the positive constant β , but never depends on $t \geq 0$. Moreover, we regard

$$\|U(t)\| = \|U(t)\|_{L^2(R^2)}, \quad \|U(t)\|_\infty = \|U(t)\|_{L^\infty(R^2)}$$

$$\|U(t)\|_m = \|U(t)\|_{H^m(R^2)}, \quad \|U_0\| = \|U_0\|_{L^2(R^2)}, \quad \|U_0\|_m = \|U_0\|_{H^m(R^2)}$$

Suppose that $f(x, y) \in L^1(R^2) \cap L^2(R^2)$, define its Fourier transform as follows

$$F[f](\zeta, \eta) = \hat{f}(\zeta, \eta) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x, y) \exp[-ix\zeta - iy\eta] dx dy$$

As usual, the definition is extended by continuity to the space of tempered distributions.