On the Generalized Porous Medium Equation in Fourier-Besov Spaces

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Abstract. We study a kind of generalized porous medium equation with fractional Laplacian and abstract pressure term. For a large class of equations corresponding to the form: $u_t + v\Lambda^{\beta}u = \nabla \cdot (u\nabla Pu)$, we get their local well-posedness in Fourier-Besov spaces for large initial data. If the initial data is small, then the solution becomes global. Furthermore, we prove a blowup criterion for the solutions.

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1 Introduction

In this paper, we study the nonlinear nonlocal equation in \mathbb{R}^n of the form

$$\begin{cases} u_t + v \Lambda^{\alpha} u = \nabla \cdot (u \nabla P u); \\ u(0, x) = u_0. \end{cases}$$
(1.1)

Usually, u = u(t,x) is a real-valued function, represents a density or concentration. The dissipative coefficient $\nu > 0$ corresponds to the viscous case, while $\nu = 0$ corresponds to the inviscid case. In this paper we study the viscid case and take $\nu = 1$ for simplicity. The fractional operator Λ^{α} is defined by Fourier transform as $(\Lambda^{\alpha} u)^{\wedge} = |\xi|^{\alpha} \hat{u}$. *P* is an abstact operator.

Equation (1.1) here comes from the same proceeding with that of the fractional porous medium equation (FPME) introduced by Caffarelli and Vázquez [5]. In fact, equation

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(1.1) comes into being by adding the fractional dissipative term $v\Lambda^{\alpha}u$ to the continuity equation $u_t + \nabla \cdot (uV) = 0$, where the velocity $V = -\nabla p$ and the velocity potential or pressure p is related to u by an abstract operator p = Pu.

The absrtact form pressure term Pu gives a good suitability in many cases. The simplest case comes from a model in groundwater in filtration [1, 20]: $u_t = \Delta u^2$, that is: v = 0, Pu = u. A more general case appears in the fractional porous medium equation [5] when v = 0 and $Pu = \Lambda^{-2s}u$, 0 < s < 1. In the critical case when s = 1, it is the mean field equation first studied by Lin and Zhang [16]. Studies on the well-posedness and regularity on those equations we refer to [4, 6, 7, 18, 19, 21, 24] and the references therein.

In the FPME, the pressure can also be represented by Riesz potential as $Pu = \Lambda^{-2s}u = \mathcal{K} * u$, with kernel $\mathcal{K} = c_{n,s}|y|^{2s-n}$. Replacing the kernel \mathcal{K} by other functions in this form: $Pu = \mathcal{K} * u$, equation (1.1) also appears in granular flow and biological swarming, named aggregation equation. The typical kernels are the Newton potential $|x|^{\gamma}$ and the exponent potential $-e^{-|x|}$.

One of concerned problems on this equation is the singularity of the potential Pu which holds the well-posedness or leads to the blowup solution. Bertozzi and Carrillo [3] show that smooth kernels at origin x = 0 lead to the global in time solution, meanwhile Li and Rodrigo [15] prove that nonsmooth kernels lead to blowup phenomenon. Li and Rodrigo [14] studied the well-posedness and blowup criterion of equation (1.1) with the pressure $Pu = \mathcal{K} * u$, where $\mathcal{K}(x) = e^{-|x|}$ in Sobolev spaces. Wu and Zhang [22] generalize their work to require $\nabla \mathcal{K} \in W^{1,1}$ which includes the case $\mathcal{K}(x) = e^{-|x|}$. They take advantage of the controllability in Besov spaces of the convolution $\mathcal{K} * u$ under this condition, as well as the controllability of its gradient $\nabla \mathcal{K} * u$.

In this article we study the well-posedness and blowup criterion of equation (1.1) in Fourier-Besov spaces under an abstract pressure condition

$$\|\widehat{\Delta_k}(\nabla P u)\|_{L^p} \le C2^{k\sigma} \|\widehat{\Delta_k} u\|_{L^p}.$$
(1.2)

In Fourier-Besov spaces, it is the localization express of the norm estimate

$$\|\nabla Pu\|_{F\dot{B}^{s}_{p,q}} \le C \|u\|_{F\dot{B}^{s+\sigma}_{p,q}}.$$
(1.3)

Corresponding to the FPME, i.e. $Pu = \Lambda^{-2s}u$, we get $\sigma = 1-2s$ obviously. And if $Pu = \mathcal{K} * u$, $\mathcal{K} \in W^{1,1}$ in the aggregation equation, we get $\sigma = 1$ when $\mathcal{K} \in L^1$ and $\sigma = 0$ when $\nabla \mathcal{K} \in L^1$.

The Fourier-Besov spaces we use here come from Konieczny and Yoneda [12] when deal with the Navier-Stokes equation (NSE) with Coriolis force. Besides, Fourier-Besov spaces have been widely used to study the well-posedness, singularity, self-similar solution, etc. of Fluid Dynamics in various of forms. For instance, the early pseudomeasure spaces PM^{α} in which Cannone and Karch studied the smooth and singular properties of Navier-Stokes equations [8]. The Lei-Lin spaces \mathcal{X}^{σ} deal with global solutions to the NSE [13] and to the quasi-geostrophic equations (QGE) [2]. The Fourier-Herz spaces \mathcal{B}_q^{σ} in the Keller-Segel system [9], in the NSE with Coriolis force [10] and in the magnetohydrodynamic equations (MHD) [17].